

Synoptic paper 2015

Phys 3017/6015

Model answers

A1

Dispersion: var. speed varies with frequency  
(or with var. length,  $\frac{dv}{d\lambda} \neq 0$ )

$$\omega = \omega(k) \quad \text{specified dispersion relation} \quad [2]$$

$$\text{Phase velocity:} \quad v = \frac{\omega}{k} \quad [1]$$

$$\text{Group velocity:} \quad g = \frac{d\omega}{dk} \quad [1]$$

A2

$$d \sin \theta_1 = n \lambda_1, \quad d \sin \theta_2 = n \lambda_2$$

$$\text{where } d = 300^{-1} 10^{-3} \text{ m} = 3.3 \cdot 10^{-6} \text{ m} \quad [1]$$

$$\text{For } n=1, \quad \Delta \theta = \arcsin \frac{\lambda_2}{d} - \arcsin \frac{\lambda_1}{d} = 0.01^\circ \quad [1]$$

$$\text{Highest order } \Rightarrow \sin \theta = 1 \Rightarrow n_{\max} = \left\lfloor \frac{d}{\lambda} \right\rfloor = \lfloor 5.7 \rfloor = 5 \quad [1]$$

$$\text{For } n=5, \quad \Delta \theta = \arcsin \frac{5\lambda_2}{d} - \arcsin \frac{5\lambda_1}{d} = 0.11^\circ \quad [1]$$

A3

Commutator

$\frac{h}{2\pi}$  Planck constant

$$[\bar{J}_x, \bar{J}_y] \equiv \bar{J}_x \bar{J}_y - \bar{J}_y \bar{J}_x = i \frac{h}{2\pi} \bar{J}_z$$

The commutator of any two components of angular momentum is  $(i\hbar) \times$  (third component)  $[\bar{J}_i, \bar{J}_j] = i \epsilon_{ijk} \bar{J}_k$  [1]

$[\bar{J}_x, \bar{J}_y] \neq 0 \Rightarrow$  there is in general no complete set of states for which both  $\bar{J}_x$  and  $\bar{J}_y$  have definite values [1]

On the other hand  $[\bar{J}^2, \bar{J}_i] = 0$  for any  $i$  [1]

$\Rightarrow$  can find complete set of states for which both  $\bar{J}^2$  and one component of  $\bar{J}$  have well-defined values. [1]

A4]

1st particle  
in state 1

2nd particle  
in state 2

2-particle wave function  $\psi(1, 2)$

Indistinguishable particles  $\Rightarrow$  exchange symmetry:

$$|\psi(2, 1)|^2 = |\psi(1, 2)|^2$$

i.e.  $\psi(2, 1) = e^{i\phi} \psi(1, 2)$

Swap twice  $\Rightarrow \psi(2, 1) = e^{2i\phi} \psi(2, 1)$

i.e.  $e^{2i\phi} = 1 \Rightarrow e^{i\phi} = \pm 1$

That is,

either  $\psi(2, 1) = \psi(1, 2)$

bosons

or

$\psi(2, 1) = -\psi(1, 2)$

fermions

integer spin  $\uparrow$   
half-integer spin  $\downarrow$

Pauli exclusion principle for fermions:

if stat. 1 = stat. 2:  $\psi(1, 1) = -\psi(1, 1) = 0$

two fermions cannot occupy same state

AS

$$\text{Ionisation energy} \sim R = 13.6 \text{ eV} \quad [1]$$

Temperature of ionisation

$$k_B T \sim R$$

[1]



$$T \sim \frac{R}{k_B} \sim \frac{13.6 \cdot 1.6 \cdot 10^{-19}}{1.38 \cdot 10^{-23}} \text{ K} \sim 1.5 \cdot 10^5 \text{ K}$$

[2]

B1

a)  $\rho = 0, \quad \underline{J} = 0 \Rightarrow \underline{\nabla} \cdot \underline{E} = 0 \quad \underline{\nabla} \cdot \underline{B} = 0$  [4]

$\underline{\nabla} \wedge \underline{E} = -\frac{\partial \underline{B}}{\partial t} \quad \underline{\nabla} \wedge \underline{B} = \mu_0 \epsilon_0 \frac{\partial \underline{E}}{\partial t}$

b)  $\underline{\nabla} \wedge (\underline{\nabla} \wedge \underline{E}) = -\frac{\partial}{\partial t} \underline{\nabla} \wedge \underline{B} \Rightarrow -\underline{\nabla}^2 \underline{E} = -\mu_0 \epsilon_0 \frac{\partial^2 \underline{E}}{\partial t^2}$  [2]

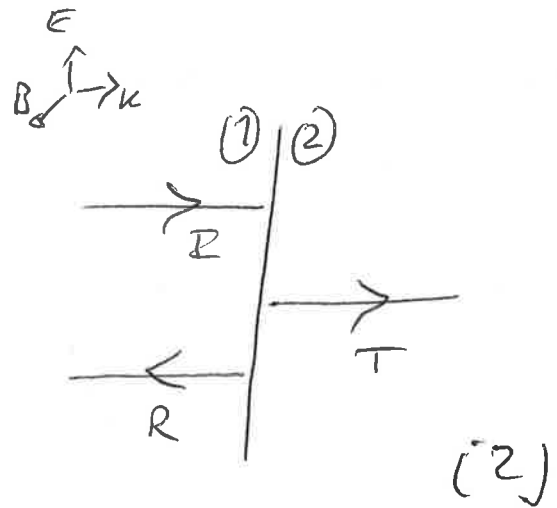
$\underline{\nabla} \underline{\nabla} \cdot \underline{E} - \underline{\nabla}^2 \underline{E} = \mu_0 \epsilon_0 \frac{\partial \underline{E}}{\partial t}$

$\underline{\nabla} \cdot \underline{E} = 0$

that is,  $\underline{\nabla}^2 \underline{E} - \frac{1}{c^2} \frac{\partial^2 \underline{E}}{\partial t^2} = 0$  [2]

wave equation with wave speed  $c = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$

c) continuity of  $\perp$  components of the fields is trivially satisfied; continuity of  $\parallel$  components:



$E_1^{\parallel} = E_2^{\parallel} \Rightarrow E_I + E_R = E_T$  [2]

$H_1^{\parallel} = H_2^{\parallel} \Rightarrow \frac{1}{v_1} E_I - \frac{1}{v_1} E_R = \frac{1}{v_2} E_T$

$\mu_1 \approx \mu_2 \approx \mu_0$

i.e.,  $(v_i \approx 1/\sqrt{\mu_0 \epsilon_i}) \quad \sqrt{\epsilon_1} E_I - \sqrt{\epsilon_1} E_R = \sqrt{\epsilon_2} E_T$  [2]

Solve for  $E_R$  and  $E_T$ :  $E_I + E_R = E_T = \sqrt{\frac{\epsilon_1}{\epsilon_2}} (E_I - E_R)$

$$\Rightarrow E_R = E_I \frac{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}}{\sqrt{\epsilon_1} + \sqrt{\epsilon_2}} \quad [2]$$

$$E_T = E_I + E_R = E_I \left( 1 - \frac{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}}{\sqrt{\epsilon_1} + \sqrt{\epsilon_2}} \right) = E_I \frac{2\sqrt{\epsilon_1}}{\sqrt{\epsilon_1} + \sqrt{\epsilon_2}} \quad [2]$$

Power per unit area:  $S = \frac{1}{\mu_0} E \frac{E}{v} = \sqrt{\frac{\epsilon}{\mu_0}} E^2 \quad [1]$

$$\Rightarrow S_I = \sqrt{\frac{\epsilon_1}{\mu_0}} E_I^2 ; S_R = \sqrt{\frac{\epsilon_1}{\mu_0}} E_R^2 ; S_T = \sqrt{\frac{\epsilon_2}{\mu_0}} E_T^2 \quad [1]$$

Energy conservation  $\Rightarrow S_I = S_R + S_T$  i.e.  $1 = \frac{S_R}{S_I} + \frac{S_T}{S_I} \quad [1]$

Verify:  $\frac{S_R}{S_I} + \frac{S_T}{S_I} = \left( \frac{E_R}{E_I} \right)^2 + \sqrt{\frac{\epsilon_2}{\epsilon_1}} \left( \frac{E_T}{E_I} \right)^2 =$

inserting expressions at the top we have

$$= \left( \frac{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}}{\sqrt{\epsilon_1} + \sqrt{\epsilon_2}} \right)^2 + \sqrt{\frac{\epsilon_2}{\epsilon_1}} \left( \frac{2\sqrt{\epsilon_1}}{\sqrt{\epsilon_1} + \sqrt{\epsilon_2}} \right)^2 =$$

$$= \frac{(\sqrt{\epsilon_1} - \sqrt{\epsilon_2})^2}{(\sqrt{\epsilon_1} + \sqrt{\epsilon_2})^2} = 1 \quad [1]$$



(B2)

a)  $H\psi = E\psi$  TISE [1]

$i\hbar \frac{\partial \psi}{\partial t} = H\psi$  TDSE [1]

The wave function of any system must satisfy the TDSE. [1]

The TISE is satisfied if the system is in a state of well-defined energy (stationary state) [1]

The wave function contains a complete set of information about the system, so it must contain complete initial conditions for solution of the time evolution equation. [1]

This would not be possible if the TDSE were second order in  $t$ . [1]

b)  $E_n = \varepsilon, 4\varepsilon, 9\varepsilon, 16\varepsilon, 25\varepsilon, \dots$

normalisation of  $\psi(0)$ :  $\|\psi(0)\|^2 = 0.2^2 + 0.3^2 + 0.4^2 + 0.843^2 = 1$

$P(E < 6\varepsilon) = P(\varepsilon \text{ or } 4\varepsilon) = 0.2^2 + 0.3^2 = 0.13$

[3]

Energy expectation value

$$\langle E \rangle = 0.2^2 \epsilon + 0.3^2 4\epsilon + 0.4^2 9\epsilon + 0.843^2 16\epsilon = 13.2\epsilon \quad [2]$$

RMS deviation

$$\langle E^2 \rangle = 0.2^2 \epsilon^2 + 0.3^2 16\epsilon^2 + 0.4^2 81\epsilon^2 + 0.843^2 256\epsilon^2 = 196.2 \epsilon^2$$

$$\Rightarrow \Delta E = \sqrt{\langle E^2 \rangle - \langle E \rangle^2} = \sqrt{196.2 \epsilon^2 - 13.2^2 \epsilon^2} = 4.69 \epsilon \quad [2]$$

$$\begin{aligned} \text{e) } \psi(t) &= 0.2 e^{-\frac{i}{\hbar} \epsilon t} \psi_1 + \\ &+ 0.3 e^{-\frac{i}{\hbar} 4\epsilon t} \psi_2 + \\ &+ 0.4 e^{-\frac{i}{\hbar} 9\epsilon t} \psi_3 + \\ &+ 0.843 e^{-\frac{i}{\hbar} 16\epsilon t} \psi_4 \end{aligned} \quad [4]$$

$$\text{d) after measurement } \psi = \psi_4 \quad [2]$$

if energy is measured again,  $E = 16\epsilon$  with certainty (1)

B3

Density of modes:

$$a) \quad g(\omega) d\omega = 2 \underset{\substack{\text{2 photon} \\ \text{polarization states}}}{g(k)} dk = 2 \frac{4\pi k^2}{(2\pi)^3} dk \quad [2]$$

$$\omega = ck \Rightarrow g(\omega) = \frac{\omega^2}{\pi^2 c^3} \quad [2]$$

$$\text{Energy density: } u = \int_0^\infty \hbar\omega \frac{1}{e^{\beta\hbar\omega} - 1} \frac{\omega^2}{\pi^2 c^3} d\omega = \quad [2]$$

$$\omega \rightarrow x = \beta\hbar\omega$$

$$= \frac{\hbar}{\pi^2 c^3} \frac{1}{(\hbar\beta)^4} \underbrace{\int_0^\infty \frac{x^3}{e^x - 1} dx}_{\pi^4/15} = \frac{\pi^2 \hbar^4 T^4}{15 \hbar^3 c^3} \quad [2]$$

$$\text{Thus } \Phi = \frac{uc}{4} = \sigma T^4 \Rightarrow \sigma = \frac{\pi^2 \hbar^4}{60 \hbar^3 c^2} \quad [2]$$

$$b) \quad \text{Total power radiated from Sun's surface} = \\ = \sigma T_s^4 4\pi R_s^2 \quad [2]$$

Isotropic emission & satellite absorbing as a black-body disk of radius  $r$

$$\Rightarrow \text{Power absorbed} = \frac{\sigma T_s^4 R_s^2}{R^2} \pi r^2 \quad [2]$$

↑  
orbit radius

In equilibrium, power absorbed = power radiated

$$\Rightarrow \frac{\sigma T_s^4 R_s^2 \pi r^2}{R^2} = \sigma T^4 4\pi r^2 \quad [2]$$

$$\text{Therefore } T = T_s \left( \frac{R_s}{2R} \right)^{1/2}$$

$$\left. \begin{array}{l} R_s = 6.955 \times 10^8 \text{ m} \\ R = 1.496 \times 10^{11} \text{ m} \\ T_s = 5700 \text{ K} \end{array} \right\} = 274.8 \text{ K} \quad [2]$$

c) With power supply on, equilibrium relation is modified to

$$W + \frac{\sigma T_s^4 R_s^2 \pi r^2}{R^2} = \sigma T^4 4\pi r^2 \quad [1]$$

$$\text{So } T = \left( \frac{W + \frac{\sigma T_s^4 R_s^2 4\pi r^2}{R^2}}{4\pi r^2 \sigma} \right)^{1/4}$$

$$= 276.5 \text{ K}$$

Now equilibrium temperature is higher by 1.7 K [1]

B4

a)  $E = T + V = \frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \dot{\phi}^2 - \frac{GM_1 m}{r}$  [1]

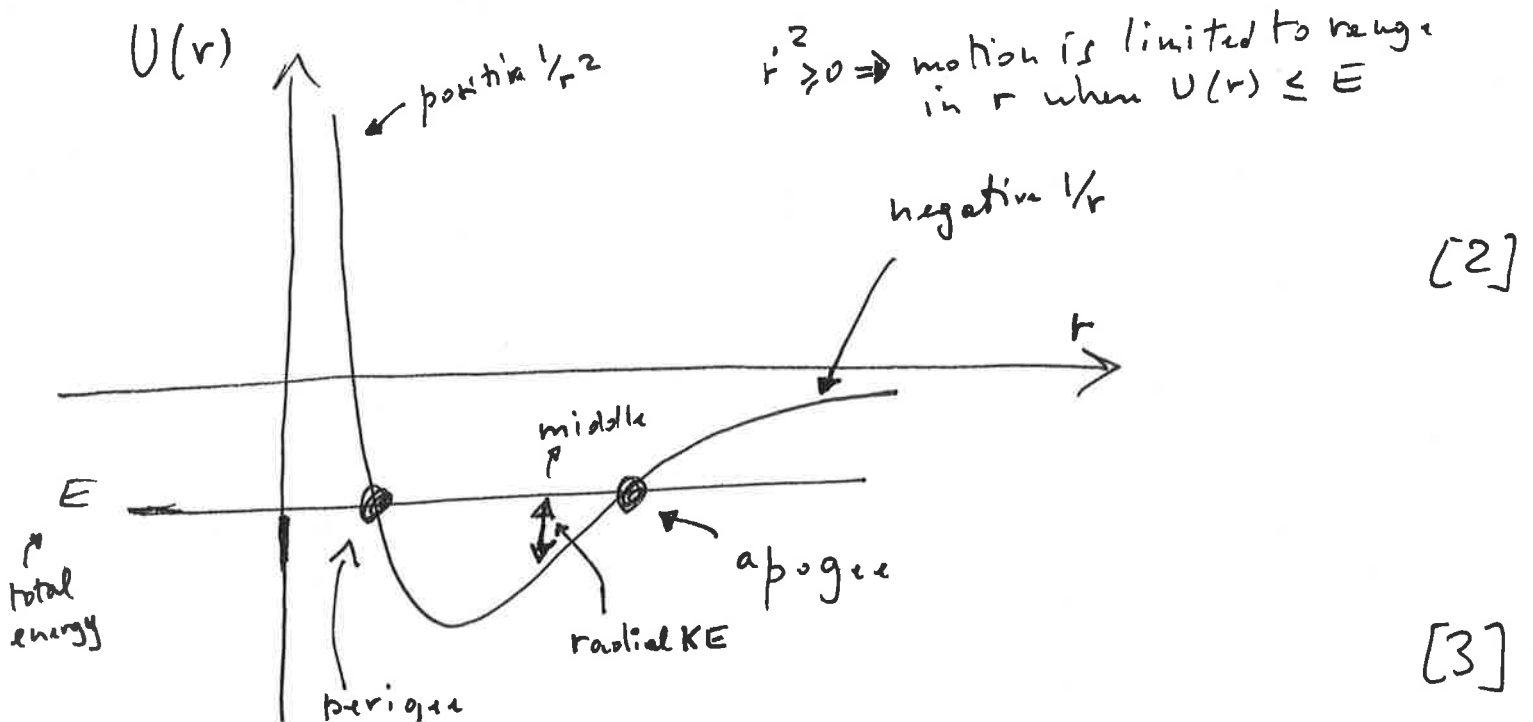
Because the gravitational force is central, angular momentum is conserved:  $J = m r^2 \dot{\phi} = \text{constant}$  [1]

$\Rightarrow E = \frac{1}{2} m \dot{r}^2 + \underbrace{\frac{J^2}{2mr^2} - \frac{GM_1 m}{r}}_{\text{effective potential } U(r)}$  "radial energy equation" [1]

radial KE  $\downarrow$     tangential KE  $\downarrow$     gravitational PE  $\downarrow$   
 $\frac{1}{2} m \dot{r}^2$      $\frac{J^2}{2mr^2}$      $-\frac{GM_1 m}{r}$

b) For given  $J$  the energy equation has the form one would get for a particle moving in one dimension in an effective potential  $U(r) = \frac{J^2}{2mr^2} - \frac{GM_1 m}{r}$  [1]

$\frac{J^2}{2mr^2}$  ← centrifugal     $-\frac{GM_1 m}{r}$  ← gravitational  
 $-\frac{\partial}{\partial r} \left( \frac{J^2}{2mr^2} \right) = \frac{J^2}{mr^3} = m r \dot{\phi}^2 = \text{centrifugal force}$  [1]

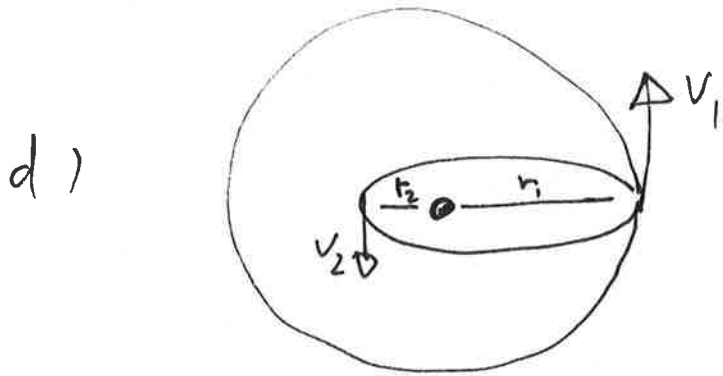


e) circular orbit:  $\frac{mv^2}{r} = \frac{GM_e m}{r^2} \Rightarrow v = \sqrt{\frac{GM_e}{r}}$  [2]

$$T = \frac{2\pi r}{v} = \frac{2\pi r^{3/2}}{\sqrt{GM_e}} \Rightarrow r = (GM_e)^{1/3} \left(\frac{T}{2\pi}\right)^{2/3}$$
 [2]

$$T = 24 \text{ h} \rightarrow \approx 4.2 \cdot 10^7 \text{ m}$$
 [1]

$$\text{Thus } v = \sqrt{\frac{GM_e}{r}} \approx 3080 \text{ m/s}$$
 [1]



Conservation of  $\vec{J} \Rightarrow m v_2 r_2 = m v_1 r_1$  [1]

Conservation of  $E \Rightarrow \frac{1}{2} m v_1^2 - \frac{GM_e m}{r_1} = \frac{1}{2} m v_2^2 - \frac{GM_e m}{r_2}$  [1]

$$\text{Then } \frac{1}{2} m v_1^2 \left[ 1 - \left(\frac{r_1}{r_2}\right)^2 \right] = GM_e m \left( \frac{1}{r_1} - \frac{1}{r_2} \right)$$

$$\Rightarrow v_1 = \sqrt{\frac{2GM_e r_2}{r_1(r_1+r_2)}} \approx 1816 \text{ m/s}$$
 [1]

So the change in velocity needed is

$$v - v_1 \approx 1260 \text{ m/s}$$
 [1]

C1

a) The spin orbit coupling comes from the interaction of the electron's magnetic moment  $\underline{\mu} = -g \frac{e}{2m} \underline{S}$  with the magnetic field seen in the rest frame of the electron as they move through the Coulomb field of the nuclear charge:  $H = -\underline{\mu} \cdot \underline{B}$ . The magnetic field is proportional to the velocity at which the electron orbits and thus to orbital angular momentum. So the interaction has the form  $-\underline{\mu} \cdot \underline{B} \propto \underline{S} \cdot \underline{L}$  [2]

As a result of spin orbit coupling  $m_l$  and  $m_s$  are no longer good quantum numbers but the total angular momentum quantum numbers  $j$  and  $m_j$  are. So a valid set is  $n, l, j, m_j$ . [1]

b)  $\underline{\mu} = -g \frac{e}{2m} \underline{S}$ ,  $g = 2$  [1]

$$\Rightarrow \underline{B} \approx -\frac{1}{mc^2} \frac{1}{r} \frac{\partial \phi}{\partial r} \underline{L}$$
$$= -\frac{1}{mc^2} \frac{1}{r} \frac{\partial}{\partial r} \left( \frac{Ze}{4\pi\epsilon_0 r} \right) \underline{L}$$
$$= \frac{1}{mc^2} \frac{Ze}{4\pi\epsilon_0 r^3} \underline{L}$$
 [1]

c)  $H(2p) \Rightarrow Z=1, n=2, l=1$  [1]

$$L \sim \hbar, \left\langle \frac{1}{r^3} \right\rangle = \left( \frac{Z}{na_0} \right)^3 \frac{2}{l(l+1)(2l+1)}$$

$$\Rightarrow B \approx \frac{e \hbar}{2m} \frac{1}{\underbrace{(4\pi \epsilon_0 c^2)}_{\frac{\mu_0}{4\pi}}} \frac{1}{a_0^3} \frac{Z^4}{\underbrace{h^3 \ell(\ell+1)(2\ell+1)}_{1/12}} \quad [1]$$

$$\approx \frac{9 \cdot 10^{-24} \cdot 10^{-7}}{(5 \cdot 10^{-11})^3 \cdot 12} \approx 0.6 \text{ T} \quad [1]$$

$$\left\langle -\frac{1}{r} \frac{\partial \phi}{\partial r} \right\rangle = \frac{Ze}{4\pi \epsilon_0} \left\langle \frac{1}{r^3} \right\rangle \Rightarrow \text{fine-structure} \sim Z^4 \quad [3]$$

$\propto Z^3$

d)  $\Delta E_n$ ; degenerate with respect to  $\ell$  and  $m_j$  [1]

For given  $j$ ,  $\ell = j - \frac{1}{2}$  or  $j + \frac{1}{2}$ , except for the highest  $j = n - \frac{1}{2}$  in which case only  $\ell = j - \frac{1}{2} = n - 1$  is possible.

For given  $j$ ,  $-j \leq m_j \leq j$  by integer steps, that is,  $2j+1$  distinct values.

Therefore the highest sublevel  $j = n - \frac{1}{2}$  is degenerate  $2j+1$  times. [1]

The other sublevels are degenerate  $2(2j+1)$  times. [1]



C2

a) Proper time = time measured by a clock in the rest frame of a moving object [2]

Proper time between two events is the interval of time between the events in a reference frame in which they occur at the same space point [1]

$$d\tau = \frac{1}{c} \sqrt{dx^\mu dx_\mu} = \frac{1}{c} \sqrt{c^2 dt^2 - dx^2}$$
$$= dt \sqrt{1 - \frac{1}{c^2} \left(\frac{dx}{dt}\right)^2} = dt \sqrt{1 - \frac{v^2}{c^2}} \quad [1]$$

$$\Rightarrow \frac{dt}{d\tau} = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} = \gamma \quad [2]$$

b) Four velocity  $u^\mu = \frac{dx^\mu}{d\tau}$

$$= \frac{dx^\mu}{dt} \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} = \left( \frac{c}{\sqrt{1 - \frac{v^2}{c^2}}}, \frac{\underline{v}}{\sqrt{1 - \frac{v^2}{c^2}}} \right) \quad [1]$$

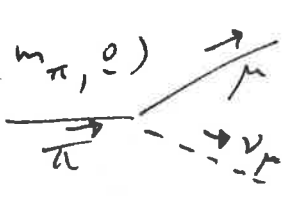
$$\Rightarrow u^\mu u_\mu = \frac{c^2}{1 - \frac{v^2}{c^2}} - \frac{v^2}{1 - \frac{v^2}{c^2}} = \frac{c^2 - v^2}{1 - \frac{v^2}{c^2}} = c^2 \quad [2]$$

$p^\mu = m u^\mu = \left( \frac{E}{c}, \underline{p} \right)$  four-momentum

$$p^\mu p_\mu = m^2 u^\mu u_\mu = m^2 c^2 \Rightarrow \quad [1]$$

$$\Rightarrow \left(\frac{E}{c}\right)^2 - \underline{p}^2 = m^2 c^2 \quad (1)$$

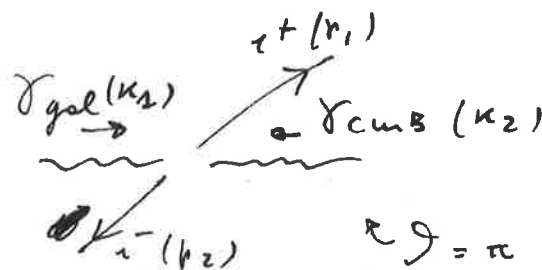
that is,  $E^2 = c^2 \underline{p}^2 + m^2 c^4$  (1)

c)  $P = (m_\pi, \underline{0})$    $p = (E, \underline{p}); p^2 = m_\mu^2$   
 $p' \quad p'^2 \approx 0$  (c=1)

4-momentum conservation  $\Rightarrow (P - p)^2 = p'^2$  (2)

i.e.  $m_\pi^2 + m_\mu^2 - 2E m_\pi = 0$

$$\Rightarrow E = \frac{m_\pi^2 + m_\mu^2}{2m_\pi} \quad \text{muon energy} \quad (2)$$

d)   $(k_1 + k_2)^2 = (p)^2$   
 $2k_1 k_2 = 2m^2 + 2p_1 p_2 \geq 4m^2$   
 $= 2 \frac{E_g E_c}{E_g E_c} (1 - \cos\theta) \leq 4 \frac{E_g E_c}{E_g E_c}$

Thus  $E_{gal} \geq \frac{4m^2}{4E_{CMB}} = \frac{m^2}{E_{CMB}}$  (2)

So the threshold energy of the galactic photons is

$$E_{\text{threshold}} = \frac{m^2}{E_{CMB}} = \frac{0.511 (10^6)^2}{2.3 \cdot 10^{-4}} \text{ eV}$$

$$= 1.14 \cdot 10^{15} \text{ eV} \quad (2)$$